

Darryl Holm's 70th Birthday Meeting:

Andrew Hone - Bi-Hamiltonian systems from cluster algebras

① History: What I worked on with Darryl

[4 papers 2001-04]

Camassa-Holm: $m_t + u m_x + 2u_x m = 0$, $m = u - u_{xx}$

• Bi-Hamiltonian: $m_t = B_j \frac{\delta H_j}{\delta m}$, $j=1,2$

$$B_1 = -(m D_x + D_x m), \quad H_1 = \int \frac{1}{2} m u dx$$

$$B_2 = D_x^3 - D_x, \quad H_2 = \int \frac{1}{2} (u u_x^2 + u^3) dx$$

Poisson brackets $\{F, G\}_j = \int \frac{\delta F}{\delta m} B_j \frac{\delta G}{\delta m} dx$

compatible: $\{, \}_1 + \{, \}_2$ satisfies Jacobi id.

• Infinite hierarchy of local symmetries (Lenard-Magri chain etc.)

$$R = B_2 B_1^{-1}; \quad m_{t_n} = R^n m_x, \quad n=0,1,2,\dots$$

Degasperis-Procesi: $m_t + u m_x + 3u_x m = 0$, $m = u - u_{xx}$

• Bi-Hamiltonian, Lax pair, conserved quantities, ..., peakons

$$u = g * m, \quad g(x) = \frac{1}{2} e^{-|x|}$$

$$m = 2 \sum_{j=1}^N p_j(t) \delta(x - q_j(t))$$

→ Hamiltonian system with N d.o.f. (non-canonical), integrable in the Liouville sense.

k-family: $m_t + u m_x + k u_x m \rightarrow$ only one Hamiltonian structure in general.

[cf. Holm-Staley]

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② outline of talk:

• Discrete integrable systems

(PΔE) Example: Hirota's discrete KdV equation

$$V_{k+1,e} - V_{k,e+1} + \alpha \left(\frac{1}{V_{k+1,e+1}} - \frac{1}{V_{k,e}} \right) = 0$$

par.
wave.

* zero curvature Lax pair: $\mathcal{L}_{k,e+1} \mathcal{M}_{k,e} = \mathcal{M}_{k+1,e} \mathcal{L}_{k,e}$

* infinitely many symmetries

Finite dimensions

(ODE) * infinitely many reductions to finite-dim.
maps: $V_{k+L,e+M} = V_{k,e} \Rightarrow V_{k,e} = V_n, n = L-kM$

$$V_{n+L+M} - V_n + \alpha \left(\frac{1}{V_{n+M}} - \frac{1}{V_{n+L}} \right) = 0 \quad (†)$$

* discrete Lax equation

→ isospectral evolution $\mathcal{L}_n^* \mathcal{M}_n^* = \mathcal{M}_n^* \mathcal{L}_{n+1}^*$

→ spectral curve $\det(\mathcal{L}_n^*(\lambda) - \mu \mathbb{1}) = 0$
generates constants of motion

Q: Birsson structure?
Lionville integrability?

A: Cluster algebras

→ Bi-Hamiltonian structure/
finite Lenard-Magri chain for (†).

③ Cluster algebras

A new class of commutative algebras, for which generators are produced by a recursive process called mutation. (Fomin & Zelevinsky 2000)

Informal definition: "A cluster algebra is a machine for generating non-trivial Laurent polynomials." (Zelevinsky)

- Appear in: Lie theory, Poisson geometry, Teichmüller theory, Solvable statistical mechanics / quantum field theory (Y-systems), discrete integrable systems, solitons, ...

Seeds: $(B, \underline{x})$, $B \in \text{Mat}_n(\mathbb{Z})$, skew $\left[\begin{array}{l} \text{coefficient-free case} \end{array} \right]$
 $\underline{x} = (x_1, \dots, x_n)$ words.
 ↑ exchange matrix ↑ cluster

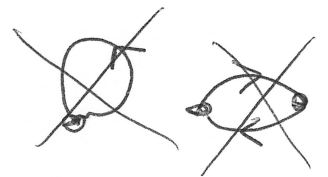
Mutation:

B

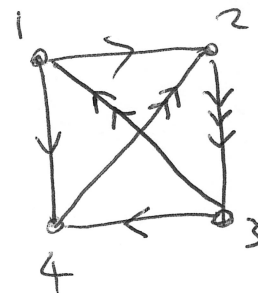
$\longleftrightarrow$

Q

quiver



e.g. $(b_{ij}) = \begin{pmatrix} 0 & 1 & -2 & 1 \\ & 0 & 3 & -2 \\ - & 0 & 1 & \\ & & & 0 \end{pmatrix} \longleftrightarrow$



Pick node k , use $[b]_+ = \max(b, 0)$
Matrix mutation:

μ_k :

$$b'_{ij} = \begin{cases} -b_{ij} & \text{if } i=k \text{ or } j=k \\ b_{ij} + [-b_{ik}]_+ b_{kj} + b_{ik} [b_{kj}]_+ & \text{otherwise.} \end{cases}$$

Quiver mutation:

① If $i \xrightarrow{p} k$ and $k \xrightarrow{q} j$
 then add $i \xrightarrow{p \times q} j$

② Reverse arrows in/out of node k .

③ Remove any 2-cycles.

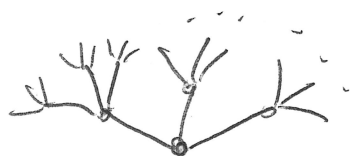
Cluster mutation:

$$\underline{\mu}_h: \quad x_j' = x_j, \quad j \neq h$$

$$x_h' = \frac{1}{x_h} \left(\prod_{i: h \rightarrow i} x_i^{[b_{ih}]_+} + \prod_{i: i \rightarrow h} x_i^{[-b_{ih}]_+} \right)$$

In general, mutation generates a tree:

e.g. $N=3$



Not a single discrete dynamical system...

Def. Cluster algebra $\mathcal{A}(B, \underline{x}) := \mathbb{Z}[\text{clusters variables}]$

Fordy-Marsh: Classification of quivers that are periodic under mutation i.e. $\mu_i(B) = \rho(B)$

$\varphi = \rho^{-1} \circ \mu$, acts trivially on B , $\uparrow$
cyclic perm

but not on $\underline{x}$ $\rightsquigarrow$ cluster map $\varphi: \underline{x} \mapsto \varphi(\underline{x})$
equivalent to recurrence

$$x_n x_{n+N} = \prod_{j=2}^{N-1} x_{n+j}^{[b_{1j+1}]_+} + \prod_{j=2}^{N-1} x_{n+j}^{[-b_{1j+1}]_+}$$

Conditions on $B = (b_{ij})$:

$$b_{1j+1} = b_{jN}$$

$$b_{j+1, n+1} = b_{jh} + b_{1j+1} [-b_{1, n+1}]_+ - b_{1, n+1} [-b_{1, j+1}]_+$$

Properties of cluster maps: ① Birational, reversible

① Volume preserving: $\Omega = \frac{dx_1 \wedge \dots \wedge dx_N}{x_1 \dots x_N}$

② Poisson brackets (sometimes trivial)

③ Presymplectic structures $\omega = \sum_{j < k} g_{jk} \frac{dx_j \wedge dx_k}{x_j x_k}$

④ Laurent property

$$x_n \in \mathbb{Z}[x_1^{\pm 1}, x_2^{\pm 1}, \dots, x_N^{\pm 1}] \quad \forall n \in \mathbb{Z}$$

⑤ Not integrable in general...

Ex. Sommes-4:

$$x_n x_{n+4} = x_{n+1} x_{n+3} + x_{n+2}^2$$

1, 1, 1, 1, 2, 3, 7, 23, 59, 314, 1529, 8209,
83313, 620297, ...

Presymplectic reduction:

$\mathbb{Z}^N = \text{im } B \oplus \ker B$
Symplectic form $\hat{\omega}$ st. $\pi^* \hat{\omega} = \omega$.
Rich basis
 $\text{im } B = \langle v_1, \dots, v_{2r} \rangle$

$$(rk B = 2r)$$

$$\pi: \underline{x} \longmapsto \underline{u} = (x^{v_1}, \dots, x^{v_{2r}})$$

$$\begin{array}{ccc} \mathbb{C}^N & \xrightarrow{\phi} & \mathbb{C}^N \\ \pi \downarrow & & \downarrow \pi \\ \mathbb{C}^{rk B} & \xrightarrow{\hat{\phi}} & \mathbb{C}^{rk B} \end{array}$$

e.g. Sommes-4: $\text{im } B = \langle (1, -2, 1, 0), (0, 1, -2, 1) \rangle$
 $\ker B = \langle (1, 1, 1, 1), (1, 2, 3, 4) \rangle$

$$(u_1, u_2) = \left(\frac{x_1 x_3}{x_2^2}, \frac{x_2 x_4}{x_3^2} \right)$$

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Symplectic map $\hat{\varphi}$, $\omega = \frac{du_1 \wedge du_2}{u_1 u_2}$

e.g. Somos-4

$$\begin{pmatrix} u_1 \\ u_2 \end{pmatrix} \mapsto$$

$$\begin{pmatrix} u_2 \\ \frac{u_2 + 1}{u_1 u_2^2} \end{pmatrix}$$

"U-system"

$$\{u_1, u_2\} = u_1 u_2$$

Conserved quantity:

$$H = u_1 u_2 + \frac{1}{u_1} + \frac{1}{u_2} + \frac{1}{u_1 u_2}$$

level sets
are
elliptic
curves

$\Rightarrow \hat{\varphi}$ is Liouville integrable (1 d.o.f.)

④ Bi-hamiltonian structures for integrable maps

(L, M) reduction of discrete KdV:

$$v_{n+L+M} - v_n + \alpha \left(\frac{1}{v_{n+M}} - \frac{1}{v_{n+L}} \right) = 0$$

(map in dim $L+M$)

Hirota tau functions:

$$v_n = \frac{\tau_n \tau_{n+L+M}}{\tau_{n+M} \tau_{n+L}}$$

$\Rightarrow$ Two bilinear eqns

period M

$$\tau_{n+2L+M} \tau_n = \beta_n \tau_{n+L+M} \tau_{n+L} + \alpha \tau_{n+2L} \tau_{n+M}$$

$$\tau_{n+2M+L} \tau_n = \beta'_n \tau_{n+L+M} \tau_{n+M} + \alpha \tau_{n+2M} \tau_{n+L}$$

period L

$\Rightarrow$ Two different cluster algebras
(with coefficients)

$\Rightarrow$ Two different presymplectic reductions

$$\dim 2 \left\lfloor \frac{L+M+1}{2} \right\rfloor - 2$$

U-system #1 : $\underline{u}$
U-system #2 : $\underline{u}'$
same dim.,
same bracket!

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e.g. $L=3, M=2$
 U -bracket has matrix entries

$$\begin{pmatrix} 0 & 0 & 1 & -1 \\ & 0 & 1 & 1 \\ & - & 0 & 0 \\ & & & 0 \end{pmatrix}$$

$$\{u_j, u_{j+1}\} = 0$$

$$\text{i.e. } \{u_j, u_{j+2}\} = u_j u_{j+2}$$

$$\{u_j, u_{j+3}\} = u_j u_{j+3}$$

same for u_j'

$$\#1: u_n u_{n+1} u_{n+2} u_{n+3} u_{n+4} = \beta_n - \alpha u_{n+2} \quad \swarrow \text{period 2}$$

$$\#2: u_n' u_{n+1}' u_{n+2}' u_{n+3}' u_{n+4}' = \beta_n' u_{n+2}' + \alpha \quad \swarrow \text{period 3}$$

dim 4
+
coeffs

KelV map $(v_0, v_1, v_2, v_3, v_4) \mapsto (v_1, v_2, v_3, v_4, v_0 + \alpha \left(\frac{1}{v_3} + \frac{1}{v_2} \right))$
 dim 5

$$\text{use } v_n = u_n u_{n+1} = u_n' u_{n+1}' u_{n+2}'$$

$\Rightarrow$ 2 different lifts of $\{, \}$ in 4D

$$\#1: \begin{aligned} \{v_0, v_1\}_1 &= v_0 v_1, \quad \{v_0, v_2\}_1 = v_0 v_2, \\ \{v_0, v_3\}_1 &= -v_0 v_3 - \alpha, \quad \{v_0, v_4\}_1 = -v_0 v_4 \end{aligned}$$

$$\#2: \begin{aligned} \{v_0, v_1\}_2 &= v_0 v_1, \quad \{v_0, v_2\}_2 = v_0 v_2 - \alpha \\ \{v_0, v_3\}_2 &= -v_0 v_3, \quad \{v_0, v_4\}_2 = -v_0 v_4 + \frac{\alpha^2}{v_2^2} \end{aligned}$$

& $\{, \}_1 + \{, \}_2$ is a Poisson bracket
 $\therefore$ compatible (bi-hamiltonian!)
 commuting first integrals I_1, I_2, I_3

Trace of monodromy $\Rightarrow$ (Lewand-Magri chain) $\uparrow \uparrow$
 Casimir of $\{, \}_2$ Casimir of $\{, \}_1$