# INVARIANT SOLUTIONS TO THE HETEROTIC EQUATIONS OF MOTION ON COMPACT QUOTIENTS OF LIE GROUPS



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#### **Objectives**

- 1. Provide new explicit smooth six-dimensional non-Kähler solutions to the Strominger system and the heterotic equations of motion.
- 2. Our solutions are based on the invariant Hermitian geometry on six-dimensional solvmanifolds studied by the authors in previous works.

## Strominger system and equations of motion

[Str86] analyzed heterotic superstring background with space-time supersymmetry. The model requires a 10-dimensional manifold of the form:

 $\mathcal{M} = \underbrace{\mathcal{M}^{3,1}}_{\text{Lorenztian}} \times \underbrace{\mathcal{M}^{6}}_{\text{Lorenztian}} , \quad g = e^{2D(x)} \begin{pmatrix} g_{\mu,\nu}(x) & 0 \\ 0 & g_{m,n}(y) \end{pmatrix}$ 

Calabi-Yau m

## **Results: Our strategy**

We choose  $(M^6, J, F, \Psi, \nabla, A)$ , where:  $\blacktriangleright M^6 = G \setminus \Gamma$  compact quotient of Lie group  $G \Longrightarrow \mathfrak{g}$  associated Lie algebra.  $\blacktriangleright$  (*J*, *F*,  $\Psi$ ) is an invariant Hermitian balanced SU(3)-structure on  $M^6$ .  $\blacktriangleright \nabla = \nabla^{\varepsilon,\rho}$ .

# $\blacktriangleright$ A is an SU(3)-instanton with respect to the previous structure.

#### **Results:** nilpotent case

- $\mathfrak{h}_3$  is the Lie algebra underlying the product of the 5-dimensional Heisenberg Lie group with  $\mathbb{R}$ . Consider: ► J:  $d\omega^1 = d\omega^2 = 0$ ,  $d\omega^3 = \omega^{1\overline{1}} - \omega^{2\overline{2}}$ .
  ► SU(3)-instance of the set o ► SU(3)-instantons:  $\nabla^+$  and  $A_\lambda$

**Equations of motion (EM)** are derived from the Lagrangian:  $S = \int e^{-2\phi} [s^{g} + 4(\nabla^{g}\phi)^{2} - \frac{1}{2}|H|^{2} - \frac{\alpha'}{4}(Tr|F^{A}|^{2} - Tr|R|^{2})]\sqrt{-g} d^{10}x,$ where  $\phi$  is a scalar field called *dilaton*. Very restrictive.

- Strominger system (SS): less restrictive than the equations of motion. If the dilaton is constant, we look for  $(M^6, J, F, \Psi, \nabla, A)$  satisfying:
- ▷ Gravitino eq.: Hol( $\nabla^+$ ) ⊆ SU(3), being  $\nabla^+$  the Bismut connection associated to the Hermitian structure (J, F). Equiv:  $\Psi$  is a holomorphic (3,0)-form i.e.  $\partial \Psi = 0$  trivializing the holomorphic canonical bundle.  $\triangleright$  Dilatino eq.: (M, J, F) is balanced, i.e.  $dF^2 = 0$ .
- $\triangleright$  Gaugino eq.: A is Donaldson-Uhlenbeck-Yau instanton, i.e.  $\Omega^A \in \mathfrak{su}(3)$ .
- The Green-Schwarz anomaly cancellation condition:

$$dT = 2\pi^2 \alpha'(\rho_1(\nabla) - \rho_1(A)) = \frac{\alpha'}{4} (\operatorname{tr}(\Omega \wedge \Omega) - \operatorname{tr}(\Omega^A \wedge \Omega^A)), \quad (1)$$

- ▶  $\alpha' \in \mathbb{R} \setminus \{0\}$  (better in physics  $\alpha' > 0$ ).
- $\blacktriangleright$   $\tau$  is the torsion 3-form associated to  $\nabla^+$ . It is identified to  $\tau = JdF$ .
- $\Omega$  is the curvature form of a metric connection  $\nabla$  on *TM*.
- Relation between (EM) and (SS) [lva10]:

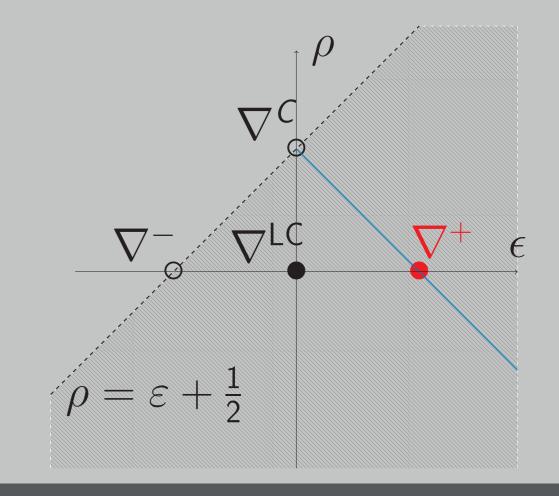
A solution to (SS) satisfies (EM)  $\iff \nabla$  in (1) is an instanton.

# $F_t = \frac{i}{2} (\omega^{1\bar{1}} + \omega^{2\bar{2}} + t^2 \omega^{3\bar{3}}). \qquad (\sigma^{A_\lambda})_2^1 = -(\sigma^{A_\lambda})_4^3 = \lambda (e^5 + e^6).$

## **Results** (with $\alpha' > 0$ ):

 $\blacktriangleright$  (SS) using  $A_{\lambda}$  with respect to  $\nabla^{\varepsilon,\rho}$  if and only if  $\rho < \varepsilon + \frac{1}{2}$ .

▶ (EM) if  $\nabla = \nabla^+$ .



#### **Results:** semisimple case

- $\mathfrak{sl}(2,\mathbb{C})$  is the Lie algebra of  $SL(2,\mathbb{C}) = \{M \in GL(2,\mathbb{C}), | \det M = 1\}$ . ► J:  $d\omega^1 = \omega^{23}$ ,  $d\omega^2 = -\omega^{13}$ ,  $d\omega^3 = \omega^{12}$ ,  $F = \frac{i}{2}(\omega^{1\bar{1}} + \omega^{2\bar{2}} + \omega^{3\bar{3}})$ .  $\blacktriangleright \nabla^{\varepsilon,\rho} \text{ is an SU(3)-instanton} \iff \nabla^{\varepsilon,\rho} = \nabla^{\mathsf{C}} \text{ (flat) or } \nabla^{\varepsilon,\rho} = \nabla^{+} \text{ (non-flat).}$ Consider:  $\beta(\varepsilon,\rho) = 1 + 4\varepsilon + 4\varepsilon^2 + 32\varepsilon^3 - 12\rho - 24\varepsilon\rho - 32\varepsilon^2\rho + 36\rho^2 + 32\varepsilon\rho^2 - 32\rho^3.$ **Results** (with  $\alpha' > 0$ ):
- ▶ If  $\beta(\varepsilon, \rho) \neq 0$ , consider  $A = \nabla^c$  and sign  $(\alpha') = \text{sign} (\beta(\varepsilon, \rho))$ . Then: (SS) for  $\nabla^{LC}$  and  $\nabla^{t<0}$  and (EM) for  $\nabla^+$ .

#### Antecedents

- There are several proposals for connections  $\nabla$  in (1):  $\triangleright$  Levi-Civita:  $\nabla^{LC}$ . Torsion-free.
  - $\triangleright$  Bismut-Strominger:  $\nabla^+ = \nabla^{LC} + \frac{1}{2}T$ , where T = JdF. ▷ Chern:  $\nabla^{C} = \nabla^{LC} + \frac{1}{2}C$ , where  $C(\cdot, \cdot, \cdot) = dF(J \cdot, \cdot, \cdot)$ .
  - $\triangleright$  Hull:  $\nabla^- = \nabla^{LC} \frac{1}{2}\overline{T}$ . ▷ Hermitian connections (Gauduchon):  $\nabla^t = \nabla^{LC} + \frac{1-t}{4}\tau + \frac{1+t}{4}C$ ,  $t \in \mathbb{R}$ .  $(\nabla^+ = \nabla^{t=-1}, \nabla^c = \nabla^{t=1}).$
- Solutions to SS & EM:
  - FIUV09 First explicit solutions based on nilmanifolds (compact quotients) of nilpotent Lie groups) with  $\alpha' > 0$ , constant dilaton and non-flat instanton A.
    - ▶ (SS) for  $\nabla = \nabla^{LC}, \nabla^+$ .
    - $\blacktriangleright$  (EM) for  $\nabla = \nabla^+$  in a nilmanifold with underlying Lie algebra  $\mathfrak{h}_3$ .
  - $\triangleright$  [UV15] Explicit solutions for (SS) based on the nilpotent Lie algebra  $\mathfrak{h}_{19}^$ with  $\alpha' > 0$ , constant dilaton, non-flat instanton and  $\nabla = \nabla^{C}$  in (1).
  - ▷ [FIUV14] non-invariant explicit solutions with non-constant dilaton for (SS) on nilmanifolds which are  $\mathbb{T}^2$ -bundles over  $\mathbb{T}^4$ .
  - $\triangleright$  [AG14], [FY15] explicit solutions for (SS) on compact quotient of  $SL(2, \mathbb{C})$ with respect to  $\nabla^{t<0}$  (flat) and  $\nabla^{t<-1}$  (non-flat).

▶ If  $\beta(\varepsilon, \rho) \neq 0$ , consider  $A = \nabla^c$  and sign  $(\alpha') = \text{sign} (\beta(\varepsilon, \rho) - 8)$ . Then, (SS) for  $\nabla^{t < -1}$ .

#### **Results: solvable case**

- $\mathfrak{g}_7$  is a solvable Lie algebra admitting lattices and complex structures with holomorphically trivial canonical bundle:
- $\blacktriangleright J: d\omega^1 = i \,\omega^1 \wedge (\omega^3 + \omega^3), \ d\omega^2 = -i \,\omega^2 \wedge (\omega^3 + \omega^3), \ d\omega^3 = \omega^{11} \omega^{22}.$ ►  $2F_u = i(\omega^{1\overline{1}} + \omega^{2\overline{2}} + \omega^{3\overline{3}}) + u\omega^{1\overline{2}} - \overline{u}\omega^{2\overline{1}}, u \in \mathbb{C}, |u| < 1.$
- ► SU(3)-instantons:  $\nabla^+$  (only if u = 0) and  $A_{\lambda,\mu}$  defined by the following connection 1-forms:  $(\sigma^{A_{\lambda,\mu}})_2^1 = -(\sigma^{A_{\lambda,\mu}})_4^3 = \lambda e^5 + \mu e^6$ .

## **Results** (with $\alpha' > 0$ ):

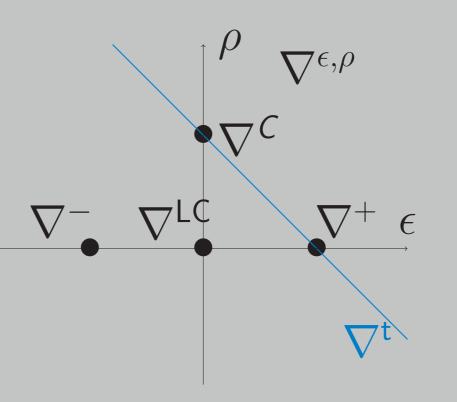
- $\blacktriangleright$  *u* = 0:
  - $\triangleright$  (SS) using  $A_{\lambda,\mu}$  with respect to  $\nabla^{\varepsilon,\rho}$  if and only if  $\rho < \varepsilon + \frac{1}{2}$ .  $\triangleright$  (EM) if  $\nabla = \nabla^+$ .
- ▶  $u \neq 0$ : (SS) using  $A_{\lambda,\mu}$  with respect to  $\nabla^{\varepsilon,\rho}$  if and only if  $(\varepsilon,\rho) \in \Delta_+$ . In particular for  $\nabla^c$ .

#### A new family of metric connections

In [OUV16] we define a new family of metric connections:

 $\nabla^{\varepsilon,\rho} = \nabla^{LC} + \varepsilon \tau + \rho c, \ (\varepsilon,\rho) \in \mathbb{R}^2.$ This family includes all the previous connections:  $abla^{LC}=
abla^{0,0},\quad 
abla^-=
abla^{-1/2,0},$  $abla^+=
abla^{1/2,0},\quad 
abla^{\mathcal{C}}=
abla^{0,1/2},$  $\nabla^t = \nabla^{\varepsilon, 1/2 - \varepsilon}.$ 

**Result:**  $\nabla^{\varepsilon,\rho}$  is Hermitian  $\iff \varepsilon + \rho = \frac{1}{2}$ .



#### References

[AG14] B. Andreas, M. García-Fernández, Commun. Math. Phys. 332 (2014), 1381–1383. [FIUV14] F. Fernández, S. Ivanov, L. Ugarte, D. Vassilev, JHEP 06 (2014) 073. [FIUV09] F. Fernández, S. Ivanov, L. Ugarte, R. Villacampa, Commun. Math. Phys. 288 (2009), 677–697. [FY15] T. Fei, S.-T. Yau, Commun. Math. Phys. 338 (2015), 1183–1195. [Iva10] S. Ivanov, Phys. Lett. B 685 (2010), 190–196. [OUV] A. Otal, L. Ugarte, R. Villacampa, arxiv. 1604.02851v1[math.DG]. [Str86] A. Strominger, Nucl. Phys. B 274 (1986), 253–284.

[UV14] L. Ugarte, R. Villacampa, Asian J. Math. 18 (2014), 229–246.

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