

Recent developments on the Strominger System

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The Strominger System

- (X, Ω) Calabi-Yau n -fold: X complex n -dimensional smooth manifold with holomorphic volume form $\Omega \in \Omega^{n,0}(X)$
- G compact semi-simple Lie group
- $P_G \rightarrow X$ principal G -bundle

Unknowns: g hermitian metric ($\omega = g(J\cdot, \cdot)$), A connection on P_G , ∇ unitary connection on (TX, g) .

Strominger System (ST)

$$F_A \wedge \omega^{n-1} = 0, \quad F_A^{0,2} = 0, \quad (1)$$

$$R_{\nabla} \wedge \omega^{n-1} = 0, \quad R_{\nabla}^{0,2} = 0, \quad (2)$$

$$d^*\omega - i(\bar{\partial} - \partial) \log \|\Omega\|_{\omega} = 0, \quad (3)$$

$$i\bar{\partial}\partial\omega - \alpha'(\text{tr } R_{\nabla} \wedge R_{\nabla} - \text{tr } F_A \wedge F_A) = 0, \quad (4)$$

Remarks: 1) typically ω non Kähler, 2) topological constraints

$$c_1(P_G) = c_1(X) = 0, \quad c_2(P_G) = c_2(X).$$

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Today, $n = 3$.

Basic example

Candelas–Horowitz–Strominger–Witten [CHSW '85] considered solutions on a Calabi-Yau 3-fold (X, Ω) with the ansatz: $P_G = (TX, g)$ and $A = \nabla$ (so $R = F$)

- Ivanov-Papadopoulos' no-go Theorem [IP '01], implies that solutions with this ansatz are precisely metrics with $SU(3)$ holonomy (X needs to be Kählerian in this case).

By Yau's solution of the Calabi Conjecture '78, such metrics are in one to one correspondence with the Kähler cone of X

$$\mathcal{K} \subset H^{1,1}(X, \mathbb{R}).$$

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A glimpse on the geometry ...

- (1) HYM: polystable holomorphic vector bundle over (X, Ω) (degree measured with respect to $\|\Omega\|_{\omega}\omega^2$).
- (2) HYM: **non-standard** holomorphic structure on TX (also polystable).
- (3) Dilatino: torsion connection with $SU(3)$ holonomy (Ricci-flat)

$$\nabla^+ = \nabla^{LC} + 1/2g^{-1}d^c\omega \quad (Bismut)$$

- (4) Bianchi identity: Flat connection on a line bundle over $C^\infty(\Sigma^2, \mathbb{C}^2 \times X)$ [Freed '86], Twisted string structures [SSS '12], Quantum sheaf cohomology [DGKS '11].

Physics: Bianchi \sim *Green-Schwarz mechanism* '84 for anomaly cancellation in heterotic string. ST \sim supersymmetric compactification with non-trivial NS 3-form flux $H = d^c\omega$ (4D, $N = 1$).

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In a nutshell.

The Strominger system provides:

- a natural generalization of the condition of $SU(3)$ -holonomy for a metric (Kähler-Ricci flat) in a non-Kähler Calabi-Yau 3-fold.
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Three problems

- *Existence problem*: characterize existence of solutions \sim Yau's Conjecture.
- *Moduli problem*: construct a moduli space and endow it with a natural geometry.
- *$(0,2)$ -mirror symmetry*: identify pairs of solutions which are physically equivalent.

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The existence problem: Yau's Conjecture

The existence problem for the Strominger system is widely open, mainly due to our lack of understanding of the Bianchi Identity.

Conjecture (Yau '06)

Let $((X, \Omega), V)$ be a pair given by a bundle V over a balanced CY3 (X, Ω, ω_0) such that

$$c_1(V) = 0, \quad c_2(V) = c_2(X) \text{ (Bott-Chern).}$$

Then, $((X, \Omega), V)$ admits a solution of the Strominger system provided that V is $[\omega_0 \wedge \omega_0]$ -polystable.

Remark: Does not specify 'where' does the metric-solution live. **General expectation:** $[\|\Omega\|_{\omega} \omega^2] = [\omega_0^2]$ (Bott-Chern).

Evidence: provided by following examples (*ad hoc* constructions, lack of general method).

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The existence problem: Kähler case

Assume (X, Ω) Kählerian.

- [Li–Yau ('05)]: Non-Kähler solutions via perturbation of [CHSW '85] solution on $TX \oplus \mathcal{O}_X$. Examples on $SU(4)$, $SU(5)$ bundles over generic quintic in \mathbb{P}^4 .
- [Andreas & GF '10]: Any stable holomorphic vector bundle V over a Kähler manifold with holonomy $= SU(3)$, such that $c_1(V) = 0$, $c_2(V) = c_2(X)$, can be perturbed to a solution of ST.

Remarks:

- AGF solves Yau's conjecture for X Kählerian. However, method has no control on $[\|\Omega\|_\omega \omega^2]$.
- Recovers Li-Yau. **Huybrechts' Theorem**: 'For a generic quintic, $TX \oplus \mathcal{O}_X$ admits stable deformations with non-trivial restriction to embedded rational curves' (\sim superpotential).
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The existence problem: non Kähler case

Some existence results:

- [Fu & Yau '08]: Goldstein–Prokushkin non Kählerian fibration (elliptic fibrations over $K3$, Monge-Ampère eq).
- [Fernandez-Ivanov-Ugarte-Villacampa '08]: Invariant solutions on nilmanifolds.
- [Fei-Yau '14]: invariant solutions on parallelizable manifolds $X = G/\Gamma$ (Wang Th.).
- [Fei-1 '15]: Singular solution on $K3$ fibration over Riemann surface.
- [Fei-2 '15]: Solution on complement of a fibre in twistor space.

Remark: for $G = SL(2, \mathbb{C})$, considered by Biswas–Mukerjee '13 and Andreas–GF '14 (parameterised by real hyperbolic 3-manifolds).

Uniqueness? How many solutions are there?

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The Moduli Problem

- *The moduli space of solutions of the Strominger system \mathcal{M}_{ST} is a fundamental gadget in string theory, where it describes basic pieces (scalar massless fields) of the 4D effective field theory induced by a $N = 1$ supersymmetric heterotic string compactification.*
- *Mathematically, \mathcal{M}_{ST} has to be constructed. Conjectured (physics prediction) to be endowed with an interesting Kähler-Hodge metric.*
- *Understanding the Kähler geometry of \mathcal{M}_{ST} should lead us to new insight for the existence and uniqueness problem (\sim Yau's Conjecture).*
- *\mathcal{M}_{ST} provides a potential ground where to study a generalization of classical mirror symmetry.*

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- *Mathematically, \mathcal{M}_{ST} has to be constructed. Conjectured (physics prediction) to be endowed with an interesting Kähler-Hodge metric.*
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The moduli problem: Elliptic curves

On an elliptic curve $E = (X, \Omega)$, the Strominger system reduces to $(\nabla = \nabla_g^{LC})$

$$F_A = 0, \quad S_g = 0.$$

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- Consider a holomorphic double extension for a suitable holomorphic Atiyah Lie algebroid V over X , constructed from a solution of ST.

$$\begin{aligned}0 &\rightarrow T^*X \rightarrow \mathcal{Q} \rightarrow V \rightarrow 0 \\0 &\rightarrow \text{ad } P^c \rightarrow V \rightarrow TX \rightarrow 0\end{aligned}$$

- Calculate infinitesimal variations of ST using the stronger anomaly cancellation condition (vs Bianchi identity) and prove they give co-cycles for the Dolbeault complex of \mathcal{Q}

$$H := d^c \omega = dB - \alpha'(CS(\nabla) - CS(A)), \quad (\text{locally})$$

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Drawback: Needs X to be a $\partial\bar{\partial}$ -manifold.

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M compact, oriented, $6d$ manifold, $L \rightarrow M$ a hermitian line bundle, \mathcal{A} space of unitary connections on L

Unknowns: $\mathcal{P} = \{(\Omega, A, \omega) : \text{satisfying (1), (2)}\} \subset \Omega^3(\mathbb{C}) \times \mathcal{A} \times \Omega^2$

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Linearisation of AST induces a complex of differential operators

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Space of infinitesimal deformations/obstructions

Theorem (____, Rubio, Tipler)

The complex is elliptic. A suitable modification of L leads to an extension of the complex S^* with identical first cohomology group

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Definition: the space of infinitesimal deformations of ST is defined as $H^1(S^*)$. The space of obstruction for ST is defined as $H^2(S^*)$.

To keep in mind! There is a well-defined map

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Suppose $p_0 = (\Omega, A, \omega) \in \mathcal{P}$ is a solution with unobstructed deformations: a neighbourhood $0 \in U \subset H^1(S^*)$ provides smooth coordinates around $[p_0]$ in the moduli space

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The exterior derivative $\delta = dflux$ defines a closed $H^3(M, \mathbb{R})$ -valued 1-form

$$\delta \in \Omega^1(\mathcal{M}_{AST}, H^3(M, \mathbb{R})),$$

and hence natural foliation on the moduli space integrating $\text{Ker } \delta$ (the *flux* only defined up to equivariant-mapping class group of L).

Anomaly cancellation and flux quantization

Want to understand the leaves of this natural foliation: try infinitesimally

$$0 \longrightarrow H^1(\mathring{S}^*) \longrightarrow H^1(S^*) \xrightarrow{\delta} H^3(X, \mathbb{R}).$$

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Down the rabbit hole ...

As notation suggests, $H^1(\mathring{S}^*)$ is the cohomology of a complex, such that

$$\mathring{S}^0 \xrightarrow{\text{here!}} S^0 \oplus \Omega^2, \quad \mathring{S}^1 = S^1 \oplus \Omega^2.$$

2-forms: play role of **symmetries**, but also as **additional data for the objects parameterised**

Substitute operator $L_4: S^1 \rightarrow \Omega^4$ by $\mathring{L}_4: \mathring{S}^1 \rightarrow \Omega^3$ (*linearized Green-Schwarz*)

$$\mathring{L}_4(\mathring{\Omega}, \mathring{a}, \mathring{\omega}, b) = db - \frac{1}{2} \left(J_0 d\mathring{\omega} - J_0 (d\mathring{\omega})^{jJ_0} + 4\alpha' \mathring{a} \wedge F \right)$$

Observation: $\mathring{P}(V, r, B) = (P(V, r), B)$ and $\mathring{L} = L_0 \oplus \dots \oplus L_3 \oplus \mathring{L}_4$ define a complex with first cohomology $H^1(\mathring{S}^*)$, provided that

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Strominger meets generalized geometry

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Defining equation for **infinitesimal symmetries of smooth transitive Courant algebroid** (Baraglia, Rubio, Hitchin) constructed from fixed solution of AST

$$E = T \oplus i\mathbb{R} \oplus T^*.$$

$$\langle V + r + \xi, W + t + \eta \rangle = \frac{1}{2}(\eta(V) + \xi(W)) - \alpha'rt$$

$$\begin{aligned} [V + r + \xi, W + t + \eta] &= [V, W] + L_V\eta - i_W d\xi + i_W i_V(d^c\omega) \\ &\quad - F(V, W) + i_V dt - i_W dr \\ &\quad - 2\alpha'(dr)t - 2\alpha'ti_VF + 2\alpha'ri_WF \end{aligned}$$

Inf. symmetries: $\text{Lie Aut } E \subset \subset \Omega^0(T \oplus i\mathbb{R}) \oplus \Omega^1(i\mathbb{R}) \oplus \Omega^2$ given by (V, r, a, B) such that $a = dr + i_VF$ and satisfy $(*)$.

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Elements in \mathring{S}^1 can be interpreted as infinitesimal variations of (in particular) **generalized metrics** $V_+ \subset E$ (here $g = \omega(\cdot, J\cdot)$)

$$V_+ = e^b \{V + r + g(X)\} \subset T \oplus i\mathbb{R} \oplus T^*.$$

Remark: a priori, no analytical tools to prove that $H^1(\mathring{S}^*)$ is finite due to defining differential equation (*) for $\text{Lie Aut } E$, but

$$0 \longrightarrow H^1(\mathring{S}^*) \longrightarrow H^1(S^*) \xrightarrow{\delta} H^3(X, \mathbb{R}).$$

Remark: the *refined space of variations* $H^1(\mathring{S}^*)$ comes closer to the physics of the heterotic string, but it is not the right space: contains a single copy of $H^{1,1}(X, \mathbb{R})$ (potentially odd dimensional).

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Complexification of the 'Kähler moduli'

Using *inner symmetries of E* , we construct an **elliptic complex \hat{S}^* of differential operators of degree 1**, such that

$$0 \rightarrow H^2(M, \mathbb{R}) \rightarrow H^1(\hat{S}^*) \rightarrow H^1(\check{S}^*) \rightarrow 0$$

We expect $H^1(\hat{S}^*)$ to be **even-dimensional**, providing natural complexification of the 'Kähler moduli' (done by hand for moduli of Calabi-Yau metrics).

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The AGSOS map

Provided that $X = (M, \Omega)$ is $\partial\bar{\partial}$ -manifold, we interpret the physical construction of Anderson–Gray–Sharpe–De la Ossa–Svanes as

$$\begin{array}{ccccccc}
 & & & & 0 & & \\
 & & & & \downarrow & & \\
 0 & \longrightarrow & H^2(M, \mathbb{R}) & \longrightarrow & H^1(\hat{S}^*) & \longrightarrow & H^1(\mathring{S}^*) \longrightarrow 0 \\
 & & & & \downarrow \text{AGSOS} & & \downarrow \\
 & & & & H^1(\mathcal{Q}) & & H^1(S^*) \\
 & & & & & & \downarrow \\
 & & & & & & H^3(M, \mathbb{R})
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The holomorphic Courant algebroid \mathcal{Q}

Observation: Every solution of the following system determines a (transitive) holomorphic Courant algebroid. In particular, every solution of the Strominger system $(\Omega, \nabla, A, \omega)$ determines a holomorphic Courant algebroid over (X, Ω) .

$$dd^c \tau^{1,1} = \alpha'(\text{tr } R \wedge R - \text{tr } F \wedge F), \quad F^{0,2} = 0, \quad R^{0,2} = 0.$$

Definition: Complexify E , and consider the integrable lift $L = e^{-i\omega} T^{0,1}$ of $T^{0,1}$. Then, Gualtieri proves that the following defines a holomorphic Courant algebroid

$$\mathcal{Q} = L^\perp / L, \quad i_{V^{0,1}} \bar{\partial}_{\mathcal{Q}} s = [e^{-i\omega} V^{0,1}, \tilde{s}] \mod L$$

(\tilde{s} lift of s to L^\perp) given by a double holomorphic extension

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$$\begin{aligned}\bar{\partial}_{\mathcal{Q}}(V + r + \xi) &= \bar{\partial}V + \bar{\partial}r + \frac{i}{2}i_VF + \bar{\partial}\xi - i(i_V\partial\omega) + \alpha' rF \\ \langle V + r + \xi, W + t + \eta \rangle &= \frac{1}{2}(\eta(V) + \xi(W)) + \alpha' rt \\ [V + r + \xi, W + t + \eta] &= [V, W] + L_V\eta - i_W d\xi \\ &\quad + [r, t] - F(V, W) + \partial_V t - \partial_W r \\ &\quad + 2\alpha'(\partial r)t + 2\alpha' t i_V F - 2\alpha' r i_W F\end{aligned}$$

Remark: Unlike $H^1(\hat{S}^*)$, $H^1(\mathcal{Q})$ has a natural complex structure

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Testing the new framework

From the point of view of physics, solutions of the Strominger system and metrics with $SU(3)$ -holonomy are two different incarnations of the same phenomenon: a $N = 1$ supersymmetric compactification of the heterotic string to 4 dimensions.

Generalized Killing spinors

M compact 6-dimensional spin manifold and E **transitive** Courant algebroid. For E obtained from reduction, $E \cong T \oplus \text{ad } P \oplus T^*$.

Generalized metric

$$E = V_+ \oplus V_-$$

Gualtieri's (generalized) connection

$$D_e^G e' = [e_-, e_+]_+ + [e_+, e_-]_- + [Ce_-, e_-]_- + [Ce_+, e_+]_+$$

where $V_+ \cong \{X + r + gX\}$ and $C(X + r + gX) = X - gX$, g Riemannian Levi-Civita connection

$$D^{LC} = D^G - \frac{1}{3} T_{D^G}$$

Given $\phi \in C^\infty(M)$, D^{LC} modified canonically to D^ϕ , comp., torsion-free.

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A unifying framework

Theorem (____, Rubio, Tipler)

On a transitive Courant algebroid obtained from reduction, the Killing spinor equations are equivalent to the Strominger system. When E is exact, a solution of the Killing spinor equations is equivalent to a metric with $SU(3)$ -holonomy. In particular, leaves of the flux foliation can be interpreted as moduli spaces of solutions of the killing spinor equations.

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Generalized geometry provides a unifying framework for the theory of the Strominger system and the well-established theory for metrics with $SU(3)$ -holonomy. Bringing in techniques from the latter, by embedding the theory into generalized geometry, is a promising approach to the existence, uniqueness and moduli problem for the Strominger system.

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Generalized geometry provides a unifying framework for the theory of the Strominger system and the well-established theory for metrics with $SU(3)$ -holonomy. Bringing in techniques from the latter, by embedding the theory into generalized geometry, is a promising approach to the existence, uniqueness and moduli problem for the Strominger system.



Thank you!



B. Andreas and M. Garcia-Fernandez, Solutions of the Strominger system via stable bundles on Calabi-Yau threefolds, *Commun. Math. Phys.* 315 (2012) 153–168.



V. Bouchard and R. Donagui, An $SU(5)$ Heterotic Standard Model, *Nucl. Phys. B* 633 (2006) 783–791.



P. Candelas, G. Horowitz, A. Strominger, E. Witten, Vacuum Configurations for Superstrings, *Nucl. Phys. B* 258 (1) (1985) 46–74.



D. Freed, Determinants, torsion and strings, *Comm. Math. Phys.* 107 (3) (1986) 483–513.



M. Fernández, S. Ivanov, L. Ugarte, R. Villacampa, Non-Kähler heterotic-string compactifications with non-zero fluxes and constant dilaton, *Commun. Math. Phys.* 288 (2009) 677–697.



J.-X. Fu and S.-T. Yau, The theory of superstring with flux on non-Kähler manifolds and the complex Monge-Ampère, *J. Diff. Geom.* 78 (2008) 369–428.



T. Fei, A Construction of Non-Kähler Calabi-Yau Manifolds and New Solutions to the Strominger System, arXiv:1507.00293 (2015).



T. Fei, Some Local Models of Heterotic Strings, arXiv:1508.05566 (2015).



T. Fei and S.-T. Yau, Invariant Solutions to the Strominger System on Complex Lie Groups and Their Quotients, arXiv:1407.7641 (2014).



M. Garcia-Fernandez, Torsion-Free Generalized Connections and Heterotic Supergravity, Commun. Math. Phys. 332 (2014) 89–115.



H. Lee, Strominger's System on non-Kähler Hermitian Manifolds, Oxford DPhil Thesis (2011).



S. Ivanov and G. Papadopoulos, Vanishing theorems and string backgrounds, Class. Quant. Grav. 18 (2001) 1089–1110, 0010038 [math.DG].



J. Li and S.-T. Yau, The existence of supersymmetric string theory with torsion, J. Diff. Geom. 70 (2005) 143-181, 0411136 [hep-th].



A. Strominger, Superstrings with torsion, Nucl. Phys. B 274 (1986) 253-284.



E. Witten, New Issues in manifolds of $SU(3)$ holonomy, Nucl. Phys. B 268 (1986) 79-112.