

Learning to predict ground state properties of gapped Hamiltonians

Provably efficient machine learning for quantum many-body problems

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joint work with **H.Y. Huang, G. Torlai, V. Albert** and **J. Preskill**

outline

- 1 motivation
- 2 proof of the main result
- 3 numerical experiments
- 4 synopsis

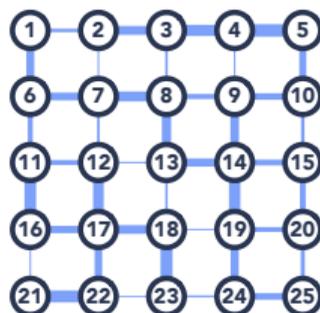
motivation

ground state problem in quantum many-body physics

Motivating numerics: 2D Heisenberg model ($n = 25, m = 40$)

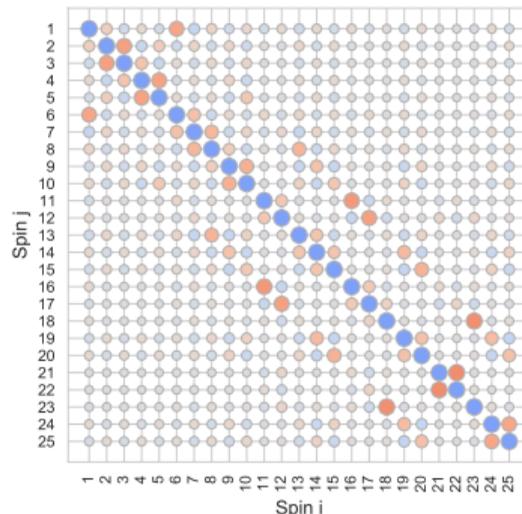
(a) 2D anti-ferromagnetic random Heisenberg model

$$H = \sum_{\langle ij \rangle} J_{ij} (X_i X_j + Y_i Y_j + Z_i Z_j)$$

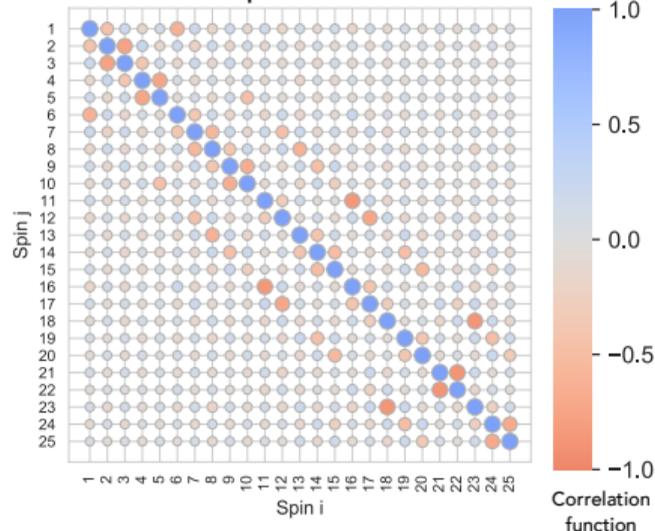


*The random J considered in (c)

(b) Exact values from DMRG



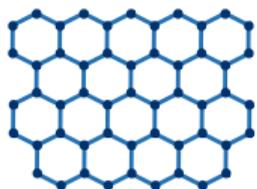
ML predictions



Correlation function

High-level vision

$$\mathbf{H}_{\text{tot}} = \sum_j \mathbf{H}$$



Parameters describing
a physical Hamiltonian



direct computation

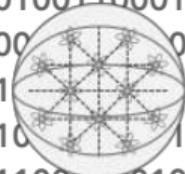
$$\rho(x) = \mathbf{v}_{\text{min}}(x) \mathbf{v}_{\text{min}}(x)^\dagger$$

(expensive: $D = 2^n$)



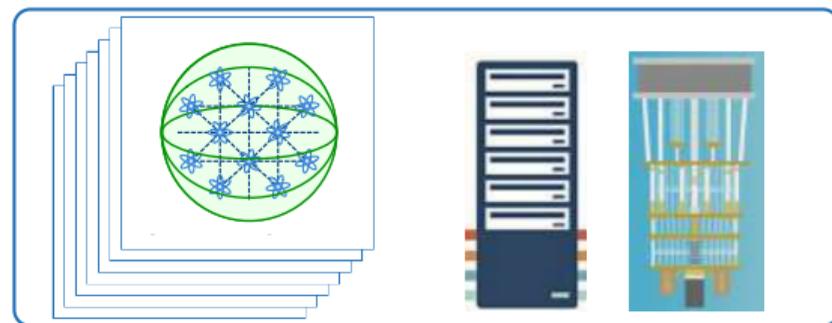
$$\text{tr}(\mathbf{O}\rho(x))$$

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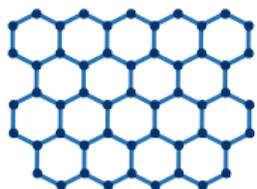
Classical representation
of the ground state

High-level vision

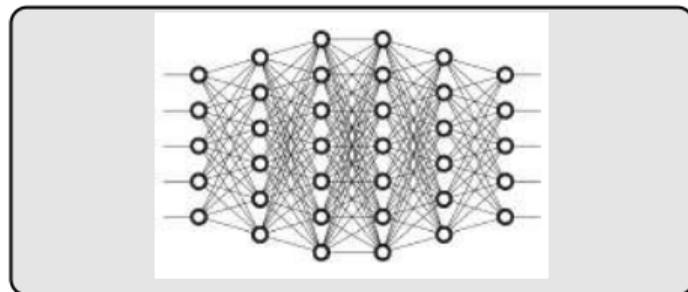


$$x \in [-1, 1]^m$$

$$\mathbf{H}_{\text{tot}}(x) = \sum_j \mathbf{H}(x)$$



 Parameters describing
a physical Hamiltonian



$$(x_\ell, \rho(x_\ell)) \Downarrow x_\ell \sim \text{unif}[-1, 1]^m$$

$$\text{tr}(\mathbf{O}\rho_{\text{train}}(x))$$

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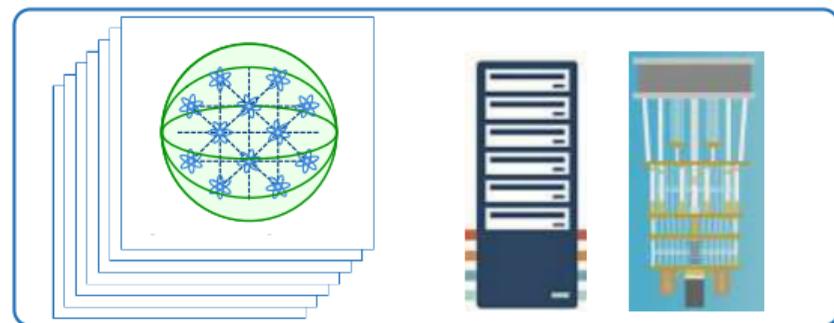
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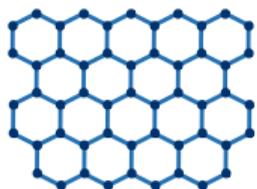
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 Parameters describing
a physical Hamiltonian



$$\rho_{\text{train}}(x) = \sum_{\ell=1}^N \kappa(x, x_{\ell}) \rho(x_{\ell})$$

κ : neural tangent kernel

$$(x_{\ell}, \rho(x_{\ell})) \Downarrow x_{\ell} \sim \text{unif}[-1, 1]^m$$



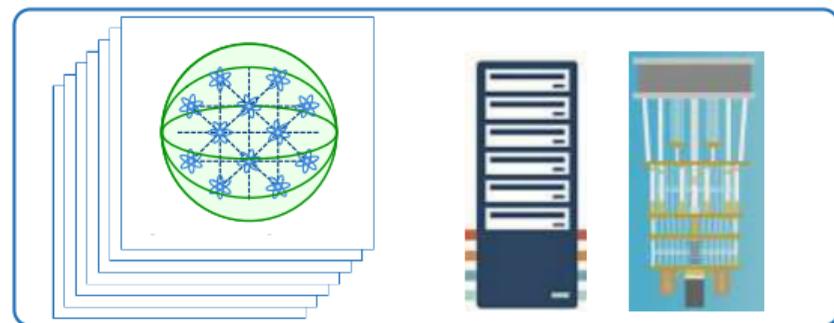
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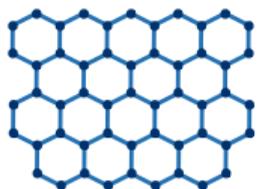
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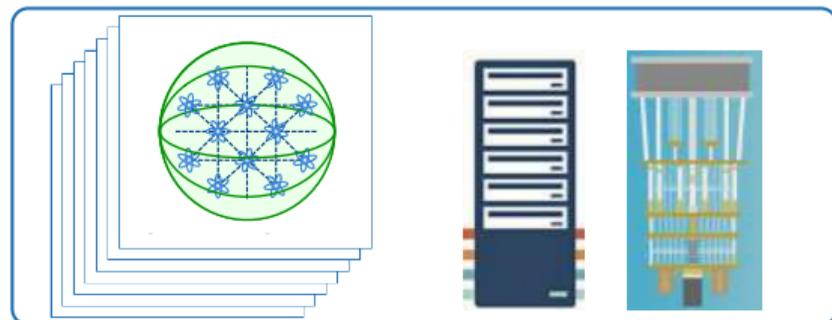
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High-level vision

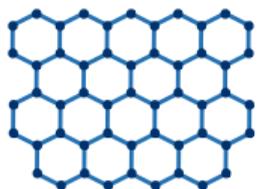
spoiler: assumptions on $H(x)$, \mathbf{O} ensure
 $\mathbb{E}_{x \sim \text{unif}[-1,1]^m} |\text{tr}(\mathbf{O}\rho_{\text{train}}(x)) - \text{tr}(\mathbf{O}\rho(x))|^2 \leq \epsilon$
 (MSE $\leq \epsilon$) with $\text{poly}(m) = \text{poly}(n)$ scaling in

- training data size
- runtime + memory



$$x \in [-1, 1]^m$$

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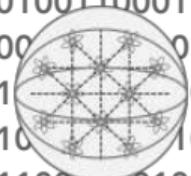
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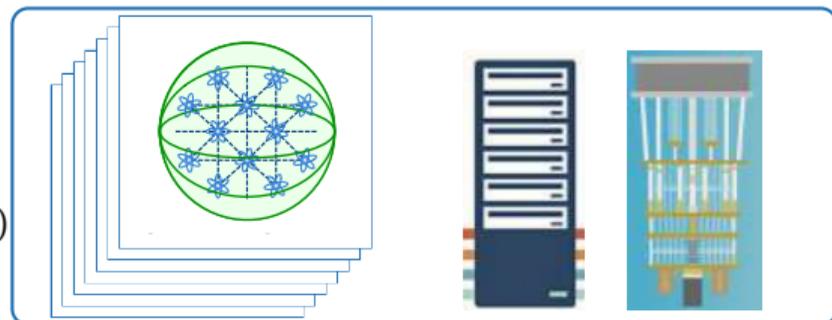
High-level vision

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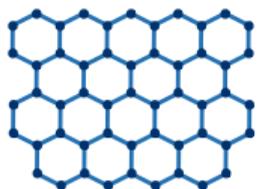
(MSE $\leq \epsilon$) with $\text{poly}(m) = \text{poly}(n)$ scaling in

- training data size (improvement to $\text{polylog}(n)$)
- runtime + memory (Lewis *et al.* 2301.13169)



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Classical representation of the ground state

Data format: classical shadows

ML insight: compression s.t. $\hat{\sigma}(\alpha\rho_1) + \beta\hat{\sigma}(\rho_2) = \hat{\sigma}(\alpha\rho_1 + \beta\rho_2)$

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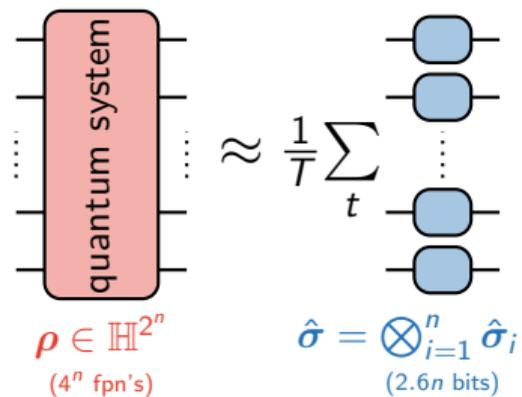
our solution : Monte Carlo sampling with quantum architectures

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our solution : Monte Carlo sampling with quantum architectures

- Monte Carlo paradigm ensures tractable approximations



$$\text{tr}(\mathbf{O}_1 \rho) \approx \frac{1}{T} \sum_{t=1}^T \text{tr}(\mathbf{O}_1 \sigma_t)$$

$$\vdots$$

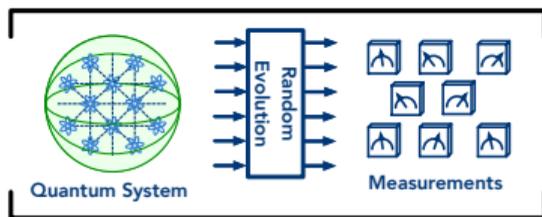
$$\text{tr}(\mathbf{O}_L \rho) \approx \frac{1}{T} \sum_{t=1}^T \text{tr}(\mathbf{O}_L \sigma_t)$$

Data format: classical shadows

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- Monte Carlo paradigm ensures tractable approximations
- sampling process outsourced to quantum simulator



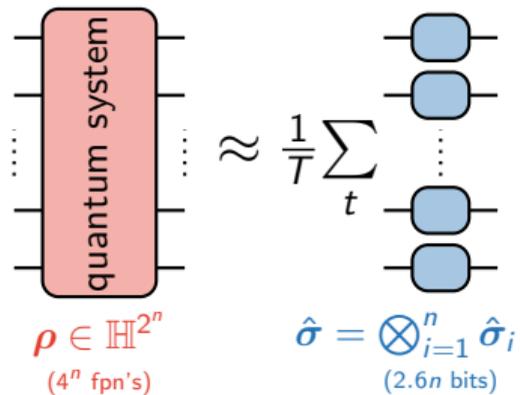
Few rounds of randomized measurements



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Classical representation
of the quantum system



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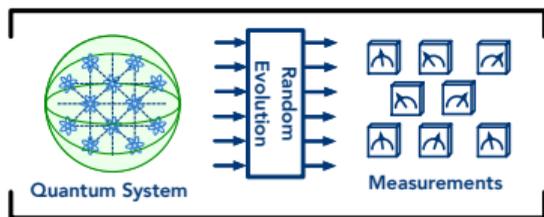
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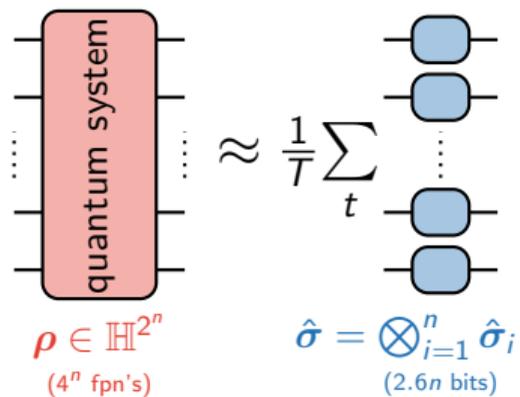


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⋮

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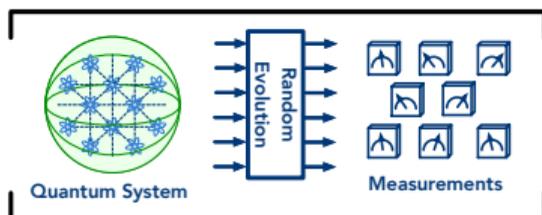
- combination of **quantum software** and **conventional software**
- simple **quantum software**, **conventional memory & runtime** is also cheap

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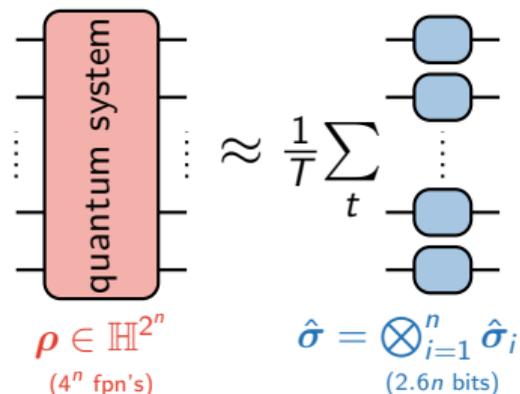
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- combination of **quantum software** and **conventional software**
- simple **quantum software**, **conventional memory & runtime** is also cheap
- works for **every state**, but only efficient for **local observables** \Rightarrow **locality assumption**

main result

Theorem 1 (Learning to predict ground state representations; informal). *For any smooth family of Hamiltonians $\{H(x) : x \in [-1, 1]^m\}$ in a finite spatial dimension with a constant spectral gap, the classical machine learning algorithm can learn to predict a classical representation of the ground state $\rho(x)$ of $H(x)$ that approximates few-body reduced density matrices up to a constant error ϵ when averaged over x . The required training data size N and computation time are polynomial in m and linear in the system size n .*

H.Y. Huang, R. Kueng, G. Torlai, V.A. Albert, J. Preskill. *Provably efficient ML for many-body problems.*

Science **377**, eabk3333 (2022) (and <https://arxiv.org/abs/2106.12627>)

proof of the main result

three steps: (i) signal processing, (ii) bridge to ground state problem, (iii) classical shadows

Proof part 1: signal processing

Theorem

Consider a function $f : [-1, 1]^m \rightarrow \mathbb{R}$ (think: $f(x) = \text{tr}(\mathbf{O}\rho(x))$) that obeys

(i) $\mathbb{E}_{x \sim \text{unif}_{[-1,1]^m}} \|\nabla_x f(x)\|_2^2 \leq C$ (controlled average gradient size)

(ii) $|f(x)| \leq B$ almost surely (bounded magnitude).

Use $N = B^2 m^{\mathcal{O}(C/\epsilon)}$ uniform samples $(x_\ell, f(x_\ell))$ with $x_\ell \stackrel{\text{unif}}{\sim} [-1, 1]^m$ to construct

$$\tilde{f}(x) = \frac{1}{N} \sum_{\ell=1}^N \kappa_\Lambda(x, x_\ell) f(x_\ell) \quad \text{with} \quad \kappa_\Lambda(x, x_\ell) = \sum_{\substack{k \in \mathbb{Z}^m \\ \|k\|_2 \leq \Lambda}} e^{i\pi \langle k, x-y \rangle}, \quad \Lambda = \mathcal{O}(C/\epsilon).$$

Then, $\mathbb{E}_{x \sim \text{unif}_{[-1,1]^m}} \left| \tilde{f}(x) - f(x) \right| \leq \epsilon$ (MSE $\leq \epsilon$) with high probability.

Proof part 1: signal processing

$$f(x) = \mathbf{F}^{-1} \mathbf{F} f(x)$$

← Fourier series plays nicely with MSE

Proof part 1: signal processing

$$\begin{aligned} f(x) &= \mathbf{F}^{-1} \mathbf{F} f(x) \\ &= \mathbf{F}^{-1} \mathbf{T}_\Lambda \mathbf{F} f(x) + \mathbf{F}^{-1} (\text{Id} - \mathbf{T}_\Lambda) \mathbf{F} f(x) \\ &\approx \mathbf{F}^{-1} \mathbf{T}_\Lambda \mathbf{F} f(x) \end{aligned}$$

← Fourier series plays nicely with MSE

← truncation in frequency domain
(use $\mathbb{E}_x \|\nabla_x f(x)\|_2^2 \leq C$)

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 &= \mathbf{F}^{-1} \mathbf{T}_\Lambda \mathbf{F} f(x) + \mathbf{F}^{-1} (\text{Id} - \mathbf{T}_\Lambda) \mathbf{F} f(x) \\
 &\approx \mathbf{F}^{-1} \mathbf{T}_\Lambda \mathbf{F} f(x) \\
 &= \sum_{\|k\|_2 \leq \Lambda} e^{i\pi \langle k, x \rangle} \frac{1}{2^m} \int_{[-1,1]^m} d^m y e^{-i\pi \langle k, y \rangle} f(y) \\
 &= \frac{1}{2^m} \int_{[-1,1]^m} \left(\sum_{\|k\|_2 \leq \Lambda} e^{i\pi \langle k, x-y \rangle} \right) f(y) d^m y
 \end{aligned}$$

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← ℓ_2 -Dirichlet kernel emerges

Proof part 1: signal processing

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 &\approx \mathbf{F}^{-1} \mathbf{T}_\Lambda \mathbf{F} f(x) && \leftarrow \text{truncation in frequency domain} \\
 &= \sum_{\|k\|_2 \leq \Lambda} e^{i\pi \langle k, x \rangle} \frac{1}{2^m} \int_{[-1,1]^m} d^m y e^{-i\pi \langle k, y \rangle} f(y) && \leftarrow \text{(use } \mathbb{E}_x \|\nabla_x f(x)\|_2^2 \leq C) \\
 &= \frac{1}{2^m} \int_{[-1,1]^m} \left(\sum_{\|k\|_2 \leq \Lambda} e^{i\pi \langle k, x-y \rangle} \right) f(y) d^m y && \leftarrow \ell_2\text{-Dirichlet kernel emerges} \\
 &= \frac{1}{2^m} \int_{[-1,1]^m} \kappa_\Lambda(x, y) f(y) d^m y \\
 &\approx \frac{1}{N} \sum_{\ell=1}^N \kappa_\Lambda(x, x_\ell) f(x_\ell) \text{ with } x_\ell \stackrel{\text{unif}}{\sim} [-1, 1]^m && \leftarrow \text{Monte Carlo approximation,}
 \end{aligned}$$

Proof part 2: bridge to ground state problem

Theorem (streamlined insight from sampling theory)

Uniform sampling efficiently interpolates functions $f : [-1, 1]^m \rightarrow \mathbb{R}$ that obey

(i) $\mathbb{E}_{x \stackrel{iid}{\sim} [-1, 1]^m} \|\nabla_x f(x)\|_2^2 \leq C$ and (ii) $|f(x)| \leq B$.

now, we set $f(x) = \text{tr}(\mathbf{O}\rho(x))$ with $\rho(x)$ ground state of $\mathbf{H}(x)$ and \mathbf{O} sum of local terms

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condition (ii): $|f(x)| = |\text{tr}(\mathbf{O}\rho(x))| \leq \|\mathbf{O}\|_\infty \|\rho(x)\|_1 = \|\mathbf{O}\|_\infty \leq B$

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condition (i) follows ground state properties of 'nice' Hamiltonians and locality
quasi-adiabatic continuation and Lieb-Robinson bounds imply

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Proposition: $\text{SPECTRALGAP}(\mathbf{H}(x)) \geq \gamma = \Omega(1)$ for all $x \in [-1, 1]^m$
ensures $\|\nabla_x \text{tr}(\mathbf{O}\rho(x))\|_2^2 \leq C_\gamma (\sum_l \|\mathbf{O}_l\|_\infty)^2 = C_\gamma B^2$ everywhere.

Proof part 2: bridge to ground state problem

Corollary (many-body restatement of main sampling theorem)

Let $\mathbf{H}(x) = \sum_j \mathbf{H}_j(x)$ with $x \in [-1, 1]^m$ be a parametrized family of 'geometrically local' n -qubit Hamiltonians with a *constant spectral gap throughout* and let $\mathbf{O} = \sum_j \mathbf{O}_j$ be a *sum of local observables* such that $\sum_j \|\mathbf{O}_j\|_\infty \leq B$. Then, a total of $N = B^2 m^{\mathcal{O}(B^2/\epsilon)}$ labeled ground states $(x_\ell, \rho(x_\ell))$ with $x_\ell \stackrel{\text{unif}}{\sim} [-1, 1]^m$ allows us to interpolate to new ground states:

$$\rho_{\text{train}}(x) = \frac{1}{N} \sum_{\ell=1}^N \kappa_\Lambda(x, x_\ell) \rho(x_\ell) \quad \text{with} \quad \kappa_\Lambda(x, x_\ell) = \sum_{k \in \mathbb{Z}^m: \|k\|_2 \leq \Lambda} e^{i\pi \langle k, x - x_\ell \rangle}, \quad \Lambda = \mathcal{O}(B^2/\epsilon).$$

With high probability, $\mathbb{E}_{x \stackrel{\text{unif}}{\sim} [-1, 1]^m} |\text{tr}(\mathbf{O}\tilde{\rho}(x)) - \text{tr}(\mathbf{O}\rho(x))|^2 \leq \epsilon$ (**MSE** $\leq \epsilon$).

- for $B = \text{const}$ and $\epsilon = \text{const}$, $N = \text{poly}(m) = \text{polylog}(D)$ (efficient training size)
- constant spectral gap is *strong physical assumption* ('deep within a phase')
- procedure is not (yet) efficient: training data $\rho(x_\ell) \in \mathbb{H}_D$ is gigantic ($D = 2^n$)

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- procedure is not (yet) efficient: training data $\rho(x_\ell) \in \mathbb{H}_D$ is gigantic ($D = 2^n$)

Proof part 3: data compression with classical shadows

- take-home message from previous slide:

$$\rho(x) \approx \rho_{\text{train}}(x) = \frac{1}{N} \sum_{\ell=1}^N \kappa_{\Lambda}(x, x_{\ell}) \rho(x_{\ell})$$

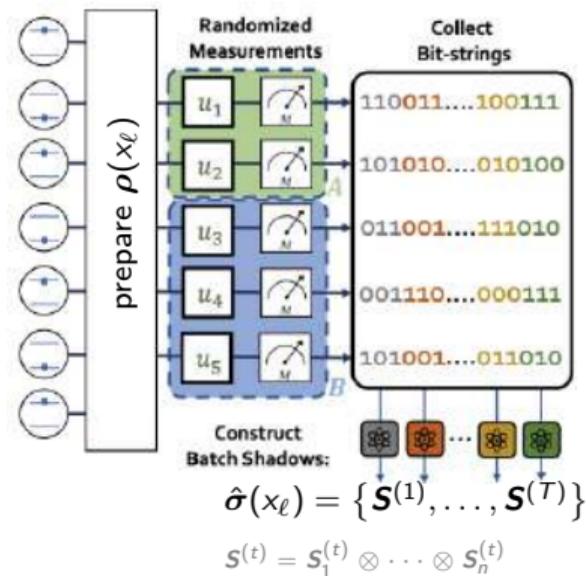
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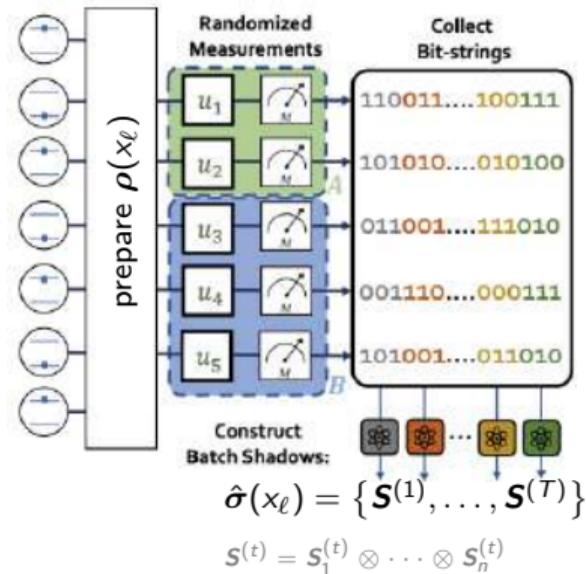
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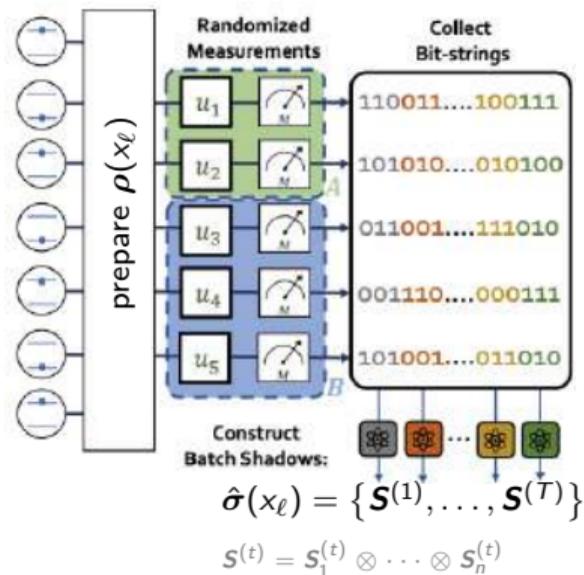
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Proposition: minimal classical shadows with $T = 1$ already ensure MSE ϵ ; moreover $\text{STORAGE}(\hat{\sigma}(x_{\ell})) = 2.6\text{bits}$
 \Rightarrow cheap data acquisition, storage + processing



Theorem (Learning to predict ground state properties with classical shadows)

Let $\mathbf{H}(x) = \sum_j \mathbf{H}_j(x)$ with $x \in [-1, 1]^m$ be a parametrized family of 'geometrically local' n -qubit Hamiltonians with a *constant spectral gap throughout* and let $\mathbf{O} = \sum_j \mathbf{O}_j$ be a *sum of local observables* such that $\sum_j \|\mathbf{O}_j\|_\infty \leq B$. Then, a total of $N = B^2 m^{\mathcal{O}(B^2/\epsilon)}$ labeled ground state sketches $(x_\ell, \sigma(x_\ell))$ (*minimal classical shadows*) with $x_\ell \stackrel{\text{unif}}{\sim} [-1, 1]^m$ allows us to interpolate to new ground state sketches:

$$\tilde{\sigma}(x) = \frac{1}{N} \sum_{\ell=1}^N \kappa_\Lambda(x, x_\ell) \sigma(\rho(x_\ell)) \quad \text{with} \quad \kappa_\Lambda(x, x_\ell) = \sum_{k \in \mathbb{Z}^m: \|k\|_2 \leq \Lambda} e^{i\pi \langle k, x - x_\ell \rangle}, \quad \Lambda = \mathcal{O}(B^2/\epsilon).$$

With high probability, this interpolation obeys $\mathbb{E}_{x \stackrel{\text{unif}}{\sim} [-1, 1]^m} |\text{tr}(\mathbf{O}\tilde{\sigma}(x)) - \text{tr}(\mathbf{O}\rho(x))|^2 \leq \epsilon$ (*MSE* $\leq \epsilon$). Moreover, all computational resources (data compression, storage, training, prediction) are bounded by $\mathcal{O}(nB^2 m^{\mathcal{O}(B^2/\epsilon)})$.

- analysis extends to infinite-width neural networks (neural tangent kernel)
- for $\epsilon, B = \text{const}$, $\mathcal{O}(nB^2 m^{\mathcal{O}(B^2/\epsilon)}) = \text{poly}(n) = \text{polylog}(D)$ (efficient cost throughout)
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Let $\mathbf{H}(x) = \sum_j \mathbf{H}_j(x)$ with $x \in [-1, 1]^m$ be a parametrized family of 'geometrically local' n -qubit Hamiltonians with a *constant spectral gap throughout* and let $\mathbf{O} = \sum_j \mathbf{O}_j$ be a *sum of local observables* such that $\sum_j \|\mathbf{O}_j\|_\infty \leq B$. Then, a total of $N = B^2 m^{\mathcal{O}(B^2/\epsilon)}$ labeled ground state sketches $(x_\ell, \sigma(x_\ell))$ (*minimal classical shadows*) with $x_\ell \stackrel{\text{unif}}{\sim} [-1, 1]^m$ allows us to interpolate to new ground state sketches:

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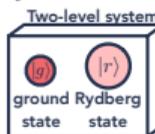
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numerical experiments

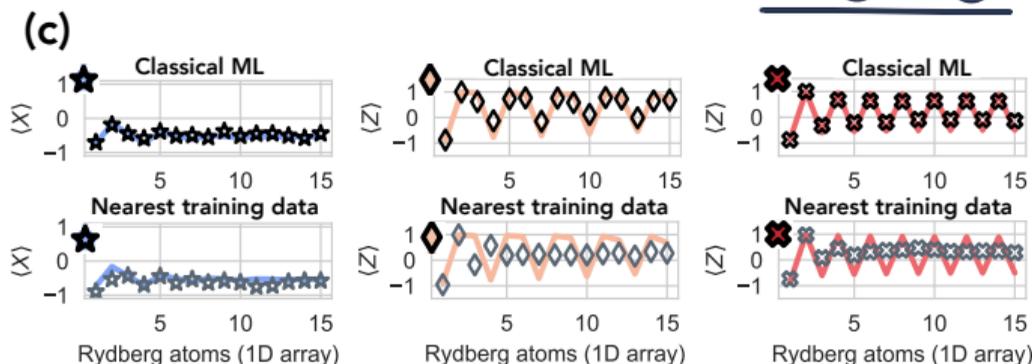
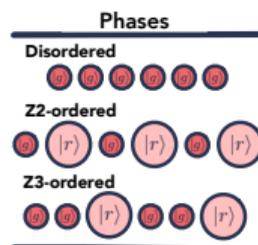
Numerics: 1D chain of $n = 51$ Rydberg atoms

(a)
$$H = \sum_i \frac{\Omega}{2} X_i - \sum_i \Delta N_i + \sum_{i < j} \Omega \left(\frac{R_b}{a|i-j|} \right)^6 N_i N_j$$

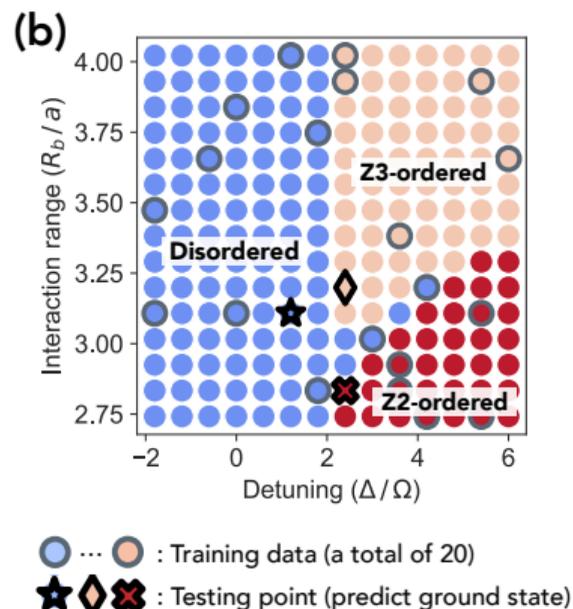
Rydberg atom array  a : atom separation

Two-level system 

$N_i = |r_i\rangle\langle r_i|$, $X_i = |g_i\rangle\langle r_i| + |r_i\rangle\langle g_i|$, $Z_i = |g_i\rangle\langle g_i| - |r_i\rangle\langle r_i|$



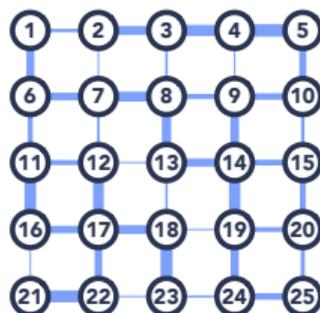
*Solid lines in the six line plots indicate exact values from DMRG



Numerics: 2D Heisenberg model with $n = 25$ spins

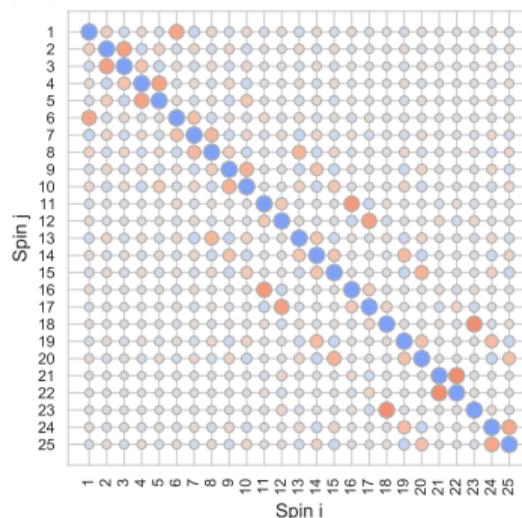
(a) 2D anti-ferromagnetic random Heisenberg model

$$H = \sum_{\langle ij \rangle} J_{ij} (X_i X_j + Y_i Y_j + Z_i Z_j)$$

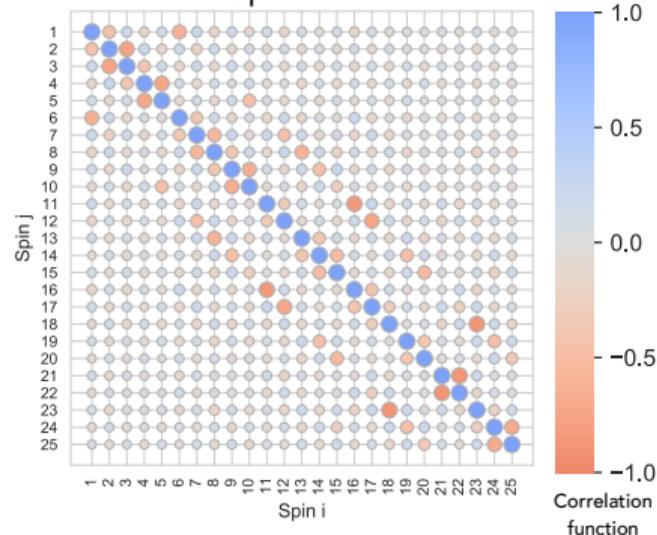


*The random J considered in (c)

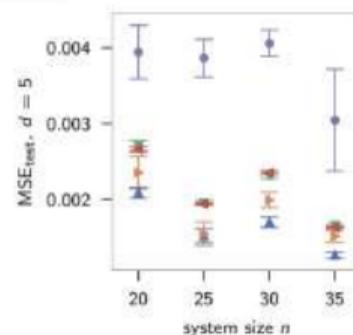
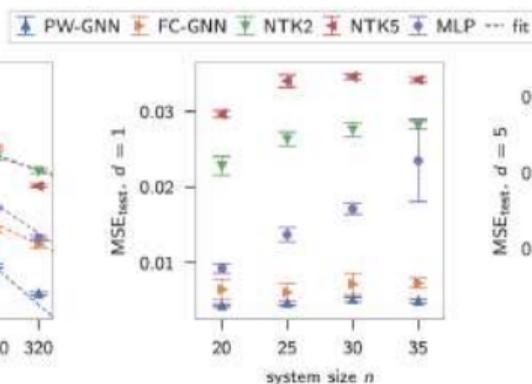
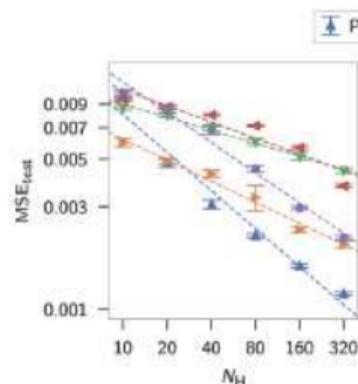
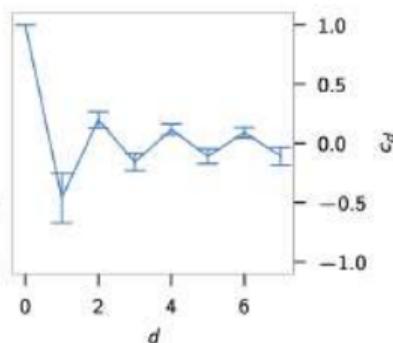
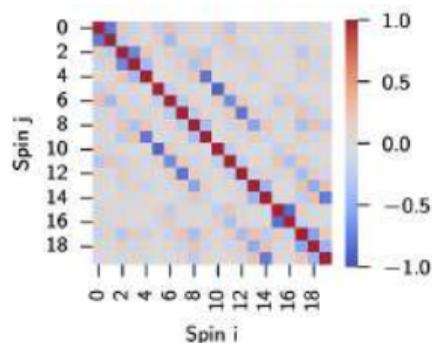
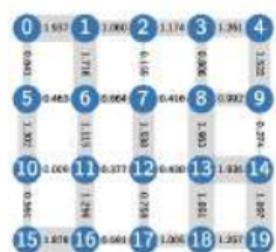
(b) Exact values from DMRG



ML predictions



Numerics: 2D Heisenberg model with different ML models



comparison between

- NKT (green, red)
- MLP (purple)
- GNN (blue, orange)

collaboration with Caltech
and Hochreiter group (JKU)
[Tran *et al*, NeurIPS workshop 22]

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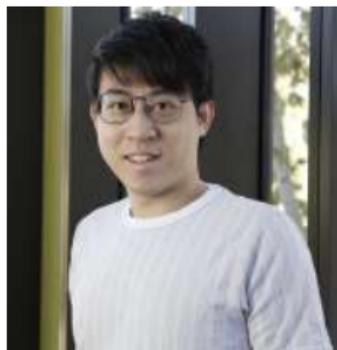
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- **complexity-theoretic bottlenecks** for approaches that don't use training data reduction from rectilinear 3-SAT and integer FACTORIZATION



H-Y. Huang



R. Kueng



G. Torlai



V. Albert

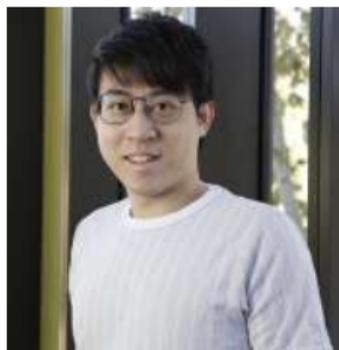


J. Preskill

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Science 377, eabk3333 (2022) (and <https://arxiv.org/abs/2106.12627>)

Thank you!



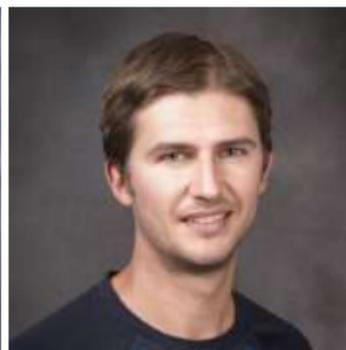
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Thank you!

Also, I am building up a team in Linz – PhD and postdoc positions are available.
Please help me spread the word.

Backup slide 1: spectral & locality control avg. gradient size

- Hamiltonian $\mathbf{H}_{\text{tot}}(x)$ has spectral gap $|\lambda_1^\uparrow(\mathbf{H}_{\text{tot}}(x)) - \lambda_2^\uparrow(\mathbf{H}_{\text{tot}}(x))| \geq \gamma = \Omega(1) \forall x \in [-1, 1]^m$
- observable $\mathbf{O}_{\text{tot}} = \sum_j \mathbf{O}_j$ is sum of 'local' terms, e.g. $\mathbf{O}_1 = \tilde{\mathbf{O}}_1 \otimes \mathbb{I}^{\otimes(n-o)}$
- **quasi-adiabatic continuation:** $W_\gamma(t)$ is fast-decaying weight function around 0 and

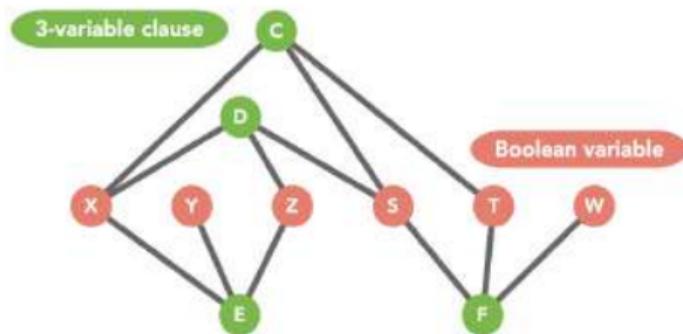
$$(\partial/\partial\hat{u})\rho(x) = [\mathbf{D}_{\hat{u}}, \rho(x)] \text{ with } \mathbf{D}_{\hat{u}} = \int_{\mathbb{R}} W_\gamma(t) e^{it\mathbf{H}_{\text{tot}}(x)} ((\partial/\partial\hat{u})\mathbf{H}_{\text{tot}}(x)) e^{-it\mathbf{H}_{\text{tot}}(x)} dt$$

$$\Rightarrow \|\nabla_x \text{tr}(\mathbf{O}\rho(x))\|_2 = \|\text{tr}([\mathbf{O}, \mathbf{D}_{\hat{u}}]\rho(x))\| \leq \|[\mathbf{O}, \mathbf{D}_{\hat{u}}]\|_\infty \|\rho(x)\|_1$$

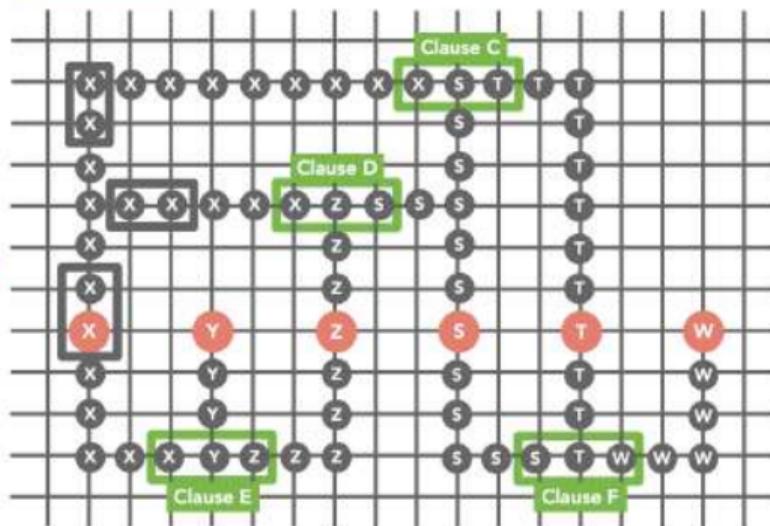
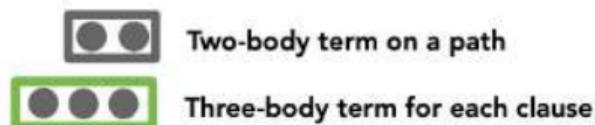
$$\leq \sum_{j_1, j_2} \int_{\mathbb{R}} W_\gamma(t) \left\| [\mathbf{O}_{j_1}, e^{it\mathbf{H}_{\text{tot}}} ((\partial/\partial\hat{u})\mathbf{H}_{j_2}) e^{-it\mathbf{H}_{\text{tot}}}] \right\|_\infty$$

- **locality** implies that most \mathbf{O}_{j_1} and \mathbf{H}_{j_2} act nontrivially on very distant regions (tensor factors)
- **Lieb-Robinson bounds** ensure that the matrix exponentials don't change this too much
- \Rightarrow almost all matrices in (1) commute approximately \Rightarrow **small gradient size** $\forall x \in [-1, 1]^m$

Backup slide 2: reduction from rectilinear 3-SAT



Planar rectilinear 3SAT Problem



Qubit Hamiltonian on 2D grid

Backup slide 3: truncation in Fourier domain

Lemma (truncation)

Let $f : [-1, 1]^m \rightarrow \mathbb{R}$ and $f_\Lambda(x) = \mathbf{F}^{-1} \mathbf{T}_\Lambda \mathbf{F} f(x) = \sum_{k \in \mathbb{Z}^m, \|k\|_2 \leq \Lambda} e^{i\pi \langle k, x \rangle} \hat{f}(k)$. Then,

$$\mathbb{E}_{x \sim \text{unif}_{[-1,1]^m}} |f(x) - f_\Lambda(x)| \leq \frac{1}{\pi^2 \Lambda^2} \mathbb{E}_{x \sim \text{unif}_{[-1,1]^m}} \|\nabla_x f(x)\|_2^2.$$

$\Rightarrow \epsilon$ -error requires cutoff $\Lambda = \mathcal{O}(C/\epsilon)$ with $C \geq \mathbb{E}_x \|\nabla_x f(x)\|_2^2$

◦ in words: average gradient size controls cutoff

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proof sketch: (standard Fourier argument with Parseval's identity)

$$\mathbb{E}_{x \sim \text{unif}_{[-1,1]^m}} \|\nabla_x f(x)\|_2^2 = \frac{1}{2^m} \int_{[-1,1]^m} d^m x \left\| \sum_{k \in \mathbb{Z}^m} i\pi k e^{i\pi \langle k, x \rangle} \hat{f}(k) \right\|_2^2 = \pi^2 \sum_{k \in \mathbb{Z}^m} \|k\|_2^2 |\hat{f}(k)|^2$$

$$\mathbb{E}_{x \sim \text{unif}_{[-1,1]^m}} |f(x) - f_\Lambda(x)| = \int_{[-1,1]^m} d^m x \left| \sum_{\|k\|_2 > \Lambda} e^{i\pi \langle k, x \rangle} \hat{f}(k) \right| = \sum_{\|k\|_2 > \Lambda} |\hat{f}(k)|^2$$

Backup slide 4: sampling approximation

Lemma (sampling rate for MSE approximation)

Let $f_\Lambda(x)$ be a band-limited function with $|f_\Lambda(x)| \leq B$ a.s. Then, with high (constant) prob:
 $\tilde{f}_\Lambda(x) = \frac{1}{N} \sum_{\ell=1}^N \kappa_\Lambda(x, x_\ell) f(x_\ell)$ with $x_\ell \stackrel{\text{unif}}{\sim} [-1, 1]^m$ and $\kappa_\Lambda(x, x_\ell) = \sum_{k \in \mathbb{Z}^m: \|k\|_2 \leq \Lambda} e^{i\pi \langle k, x - x_\ell \rangle}$
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$$a = \mathbb{E}_{x \stackrel{\text{unif}}{\sim} [-1, 1]^m} \left| \tilde{f}_\Lambda(x) - f_\Lambda(x) \right|^2 \leq B^2 m^{\mathcal{O}(\Lambda^2)} / N.$$

in words: bounded function value implies bounded (uniform) sampling rate N

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proof sketch: (polynomial bound on lattice size + Hoeffding)

$$\begin{aligned} a &= \frac{1}{2^m} \int_{[-1, 1]^m} d^m x \left| \sum_{\|k\|_2 \leq \Lambda} e^{i\pi \langle k, x \rangle} \left(\frac{1}{N} \sum_{\ell=1}^N (e^{-i\pi \langle k, x_\ell \rangle} f(x_\ell) - \mathbb{E}_{x_\ell} [e^{-i\pi \langle k, x_\ell \rangle} f(x_\ell)]) \right) \right|^2 \\ &= \sum_{k \in \mathbb{Z}^m: \|k\|_2 \leq \Lambda} \left| \frac{1}{N} \sum_{\ell=1}^N (e^{-i\pi \langle k, x_\ell \rangle} f(x_\ell) - \mathbb{E}_{x_\ell} [e^{-i\pi \langle k, x_\ell \rangle} f(x_\ell)]) \right|^2 \end{aligned}$$

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